

#### Connecting orbits and invariant manifolds in the spatial three-body problem

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### Introduction

#### Goal

□ Use dynamical systems techniques to identify key transport mechanisms and useful orbits for space missions.

#### Outline

Circular restricted three-body problem

- Equilibrium points and invariant manifold structures
- Construction of trajectories with prescribed itineraries
- Connecting orbits, e.g., heteroclinic connections
- □ Tours of Jupiter's moons

#### Introduction

#### Current research importance

Development of some NASA mission trajectories, such as the recently launched Genesis Discovery Mission, and the upcoming Jupiter Icy Moons Orbiter
 Of current astrophysical interest for understanding the transport of solar system material (eg, how ejecta from Mars gets to Earth, etc.)

#### Circular restricted 3-body problem

- the two primary bodies move in circles; the much smaller third body moves in the gravitational field of the primaries, without affecting them
- the two primaries could be Jupiter and a moon
- □ the smaller body could be a spacecraft or asteroid
- we consider the planar and spatial problems
- □ there are five equilibrium points in the rotating frame, places of balance which generate interesting dynamics

#### Circular restricted 3-body problem

Consider the two unstable points on line joining the two main bodies –  $L_1, L_2$ 



Equilibrium points –  $L_1, L_2$ 

- $\Box$  orbits exist around  $L_1$  and  $L_2$ ; both periodic and quasiperiodic
  - Lyapunov, halo and Lissajous orbits

one can draw the invariant manifolds assoicated to  $L_1$  (and  $L_2$ ) and the orbits surrounding them

these invariant manifolds play a key role in what follows

Equations of motion (planar):

$$\ddot{x} - 2\dot{y} = -\bar{U}_x, \quad \ddot{y} + 2\dot{x} = -\bar{U}_y$$

where

$$\bar{U} = -\frac{(x^2 + y^2)}{2} - \frac{1 - \mu}{r_1} - \frac{\mu}{r_2}.$$

□ Have a first integral, the Hamiltonian energy, given by  $E(x,y,\dot{x},\dot{y}) = \frac{1}{2}(\dot{x}^2 + \dot{y}^2) + \bar{U}(x,y).$ 

Energy manifolds are codimension 1 in the phase space.

### **Realms of Possible Motion**

#### **Effective potential**

□ In a rotating frame, the equations of motion describe a particle moving in an effective potential plus a magnetic field (goes back to work of Jacobi, Hill, etc).

#### **Realms of Possible Motion**



Effective potential

Level set shows accessible realms

### Motion Near Equilibria

#### For saddles of rank 1

- □ Near equilibrium point, suppose linearized Hamiltonian vector field has eigenvalues  $\pm i\omega_j$ , j = 1, ..., N - 1, and  $\pm \lambda$ .
- Assume the complexification is diagonalizable.
   Hamiltonian normal form theory tranforms Hamiltonian into a lowest order form:

$$H(q, p) = \sum_{i=1}^{N-1} \frac{\omega_i}{2} \left( p_i^2 + q_i^2 \right) + \lambda q_N p_N.$$

□ Equilibrium point is of type center  $\times \cdots \times$  center  $\times$  saddle (N - 1 centers).

#### Motion Near Equilibria

#### Multidimensional "saddle point"

For fixed energy H = h, energy surface  $\simeq S^{2N-2} \times \mathbb{R}$ .
Constants of motion:  $I_j = q_j^2 + p_j^2, j = 1, \dots, N-1$ , and  $I_N = q_N p_N$ .



The N Canonical Planes

#### Motion Near Equilibria

□ Normally hyperbolic invariant manifold at  $q_N = p_N = 0$ , n-1

$$\mathcal{M}_{h} = \sum_{i=1}^{n-1} \frac{\omega_{i}}{2} \left( p_{i}^{2} + q_{i}^{2} \right) = h > 0.$$

Note that  $\mathcal{M}_h \simeq S^{2N-3}$ , not a single trajectory.

□ Four "cylinders" of asymptotic orbits: the stable and unstable manifolds  $W^s_{\pm}(\mathcal{M}_h), W^u_{\pm}(\mathcal{M}_h)$ , which have the structure  $S^{2N-3} \times \mathbb{R}$ .

#### □ Dynamics near $L_1 \& L_2$ in spatial problem: saddle × center × center.

 $\Box \text{ Hamiltonian for linearized equations has eigenvalues} \\ \pm \lambda, \pm i\nu, \text{ and } \pm i\omega, \text{ where } \nu \neq \omega,$ 

□ Change of coordinates yields

$$H_2 = \lambda q_1 p_1 + \frac{\nu}{2} (q_2^2 + p_2^2) + \frac{\omega}{2} (q_3^2 + p_3^2).$$

For fixed energy H = h, energy surface  $\simeq S^4 \times \mathbb{R}$ .
Constants of motion:  $q_1p_1$ ,  $q_2^2 + p_2^2$  and  $q_3^2 + p_3^2$ .

Normally hyperbolic invariant manifold at  $q_1 = p_1 = 0$ ,  $\mathcal{M}_h = \frac{\nu}{2}(q_2^2 + p_2^2) + \frac{\omega}{2}(q_3^2 + p_3^2) = h > 0.$ Note that  $\mathcal{M}_h \simeq S^3$ , not a single trajectory. □ Four "cylinders" of asymptotic orbits: the stable and unstable manifolds  $W^s_+(\mathcal{M}_h), W^u_+(\mathcal{M}_h)$ , which have the structure  $S^3 \times \mathbb{R}$ .

- **B** : **bounded orbits** (periodic/quasi-periodic):  $S^3$  (3-sphere)
- A : asymptotic orbits to 3-sphere:  $S^3 \times I$  ("tubes")
- T : transit and NT : non-transit orbits.



The flow in the equilibrium region.

- **B** : **bounded orbits** (periodic/quasi-periodic): S<sup>3</sup> (3-sphere)
- A : asymptotic orbits to 3-sphere:  $S^3 \times I$  ("tubes")
- T : transit and NT : non-transit orbits.



Projection to configuration space.

Asymptotic orbits form **4D invariant manifold tubes**  $(S^3 \times I)$  in 5D energy surface.

 $\Box$  red = unstable, green = stable



- These manifold tubes play an important role in governing what orbits approach or depart from a moon (**transit orbits**)
- and orbits which do not (**non-transit orbits**)
- transit possible for objects "inside" the tube, otherwise no transit — this is important for transport issues



 Transit orbits can be found using a Poincaré section transversal to a tube.



- One can systematically construct new trajectories, which use little fuel.
  - by linking stable and unstable manifold tubes in the right order
  - and using Poincaré sections to find trajectories "inside" the tubes
- One can construct trajectories involving multiple 3-body systems.

- For a single 3-body system, we wish to link invariant manifold tubes to construct an orbit with a desired itinerary
- Construction of (X; M, I) orbit.



The tubes connecting the X, M, and I regions.

- First, integrate two tubes until they pierce a common Poincaré section transversal to both tubes.
- Second, pick a point in the region of intersection and integrate it forward and backward.

- Integrate two tubes
- Integrate a point in the region of intersection



- **Planar**: tubes  $(S \times I)$  separate transit/non-transit orbits.
- Red curve  $(S^1)$  : slice of  $L_2$  unstable manifold Green curve  $(S^1)$  : slice of  $L_1$  stable manifold
- Any point inside the intersection region  $\Delta_M$  is a (X; M, I) orbit.



- **Spatial**: Invariant manifold tubes  $(S^3 \times I)$
- Poincaré slice is a topological 3-sphere S<sup>3</sup> in ℝ<sup>4</sup>.
   S<sup>3</sup> looks like disk × disk: ξ<sup>2</sup> + ξ<sup>2</sup> + η<sup>2</sup> + ή<sup>2</sup> = r<sup>2</sup> = r<sup>2</sup><sub>ξ</sub> + r<sup>2</sup><sub>η</sub>



- Similarly, while the cut of the stable manifold tube is  $S^3$ , its projection on  $(y, \dot{y})$  plane is a curve for  $z = c, \dot{z} = 0$ .
- Any point inside this curve is a (M, I) orbit.
- Hence, any point inside the intersection region  $\Delta_M$  is a (X;M,I) orbit.





## **Connecting Orbits**



An  $L_1$ - $L_2$  heteroclinic connection

#### Tours of Jupiter's Moons

# **Tours of planetary satellite systems.** □ Example 1: Europa → Io → Jupiter



#### **Tours of Jupiter's Moons**

# $\Box Example \ 2$ : Ganymede $\rightarrow$ Europa $\rightarrow$ injection into Europa orbit



### **Tours of Jupiter's Moons**

#### • The **Petit Grand Tour** can be constructed as follows:

- Approximate 4-body system as 2 nested **3-body systems**.
- Choose an appropriate Poinaré section.
- Link the invariant manifold tubes in the proper order.
- Integrate initial condition (patch point) in the 4-body model.



#### Some References

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