### **Transport in Hamiltonian Systems With Two or More Degrees of Freedom** *Shane Ross*

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**Control and Dynamical Systems** 

# Outline

### Transport theory

 $\Box$  Time-independent Hamiltonian systems

 $\Box$  with 2 degrees of freedom

- $\Box$  with 3 (or N) degrees of freedom
  - Example: restricted three-body problem

## **Chaotic Dynamics**



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# **Transport Theory**

- Chaotic dynamics  $\rightarrow$  statistical methods
- Transport theory
  - $\Box$  Motion of ensembles of trajectories in phase space
  - $\Box$  Asks: How long to move from one region to another?
  - Determine transition probabilities, correlation functions
  - $\Box$  Applications:
    - Atomic ionization rates
    - Chemical reaction rates
    - Comet transition rates
    - Asteroid collision probabilities



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### Systems with potential barriers

• Electron near a nucleus



#### • Comet near the Sun and Jupiter



### Partition is specific to problem

□ We desire a way of describing dynamical boundaries that represent the "frontier" between qualitatively different types of behavior

### **Example:** motion of comet

- $\Box$  motion around Sun
- $\Box$  motion around Jupiter

# **Statement of Problem**

 $\Box$  Suppose we study the motion on a manifold  $\mathcal{M}$  $\Box$  Suppose  $\mathcal{M}$  is partitioned into disjoint regions

$$R_i, i=1,\ldots,N_R,$$

such that

$$\mathcal{M} = igcup_{i=1}^{N_R} R_i.$$

- □ To keep track of the initial condition of a point, we say that *initially* (at t = 0) region  $R_i$  is uniformly covered with species  $S_i$ .
- $\Box$  Thus, species type of a point indicates the region in which it was located initially.

## **Statement of Problem**

Statement of the transport problem: **Describe the distribution of species**   $S_i, i = 1, ..., N_R$ , throughout the regions  $R_j, j = 1, ..., N_R$ , for any time t > 0.



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# **Statement of Problem**

#### $\Box$ Some quantities we would like to compute are:

- $T_{i,j}(t)$  = the total amount of species  $S_i$  contained in region  $R_j$  at time t
- $F_{i,j}(t) = \frac{dT_{i,j}}{dt}(t) =$  the flux of species  $S_i$  into region  $R_j$  at time t



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# Hamiltonian Systems

### **Time-independent Hamiltonian** H(q, p)

- $\Box N$  degrees of freedom
- $\Box$  Motion constrained to a (2N 1)-dimensional energy surface  $\mathcal{M}_E$  corresponding to a value H(q, p) = E = constant

 $\Box$  Symplectic area is conserved along the flow

$$\oint_{\mathcal{L}} p \cdot dq = \int_{\mathcal{A}} dp \wedge dq = \text{constant}$$

# Symplectic Area Conserved

$$\sum_{i=1}^{N} \sigma_i \int_{\mathcal{A}^i} dp_i dq^i = \text{constant on an energy surface}$$



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## **Poincaré Section**

□ Suppose there is another (2N - 1)-dimensional surface Q that is transverse (i.e., nowhere parallel) to the flow in some local region.
□ The Poincaré section S is the (2N - 2)-dimensional intersection of M<sub>E</sub> with Q.



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## **Example for** N = 2

Circular restricted 3-body prob. (2D)

$$H = \frac{1}{2}((p_x + y)^2 + (p_y - x)^2) + U^{\text{eff}}(x, y)$$



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# 3-Body Problem (2D)

#### Look at fixed energy



#### **Position Space Projections**

# 3-Body Problem (2D)

#### **Partition the energy surface**



#### **Position Space Projection**

# 3-Body Problem (2D)

#### Look at motion near "saddle points"



**Position Space Projection** 

## **Potential Barriers**

□ Hamiltonian systems with potential barriers give rise to "saddle points" whose local form is given by

$$H(q,p) = \frac{\omega}{2}(q_1^2 + p_1^2) + \lambda q_2 p_2, \qquad (1)$$

i.e., linearized vector field has eigenvalues  $\pm i\omega$ ,  $\pm\lambda$ .

- □ Moser [1958] showed that the qualitative behavior of
   (1) carries over to the full nonlinear equations.
- $\Box$  In particular, the flow of (1) has form center  $\times$  saddle.

## Local Dynamics

 $\Box$  For fixed energy H = h, energy surface  $\simeq S^2 \times \mathbb{R}$ .  $\Box$  Other constants of motion:  $I_1 = q_1^2 + p_1^2$  and  $I_2 = q_2 p_2$ .



 $\Box$  Normally hyperbolic invariant manifold at  $q_2 = p_2 = 0$ , i.e.,

$$\mathcal{M}_h = \frac{\omega}{2}(q_1^2 + p_1^2) = h > 0.$$
  
Note that  $\mathcal{M}_h \simeq S^1$ , a periodic orbit.

# Local Dynamics

Four cylinders of asymptotic orbits: the stable and unstable manifolds  $W^s_{\pm}(\mathcal{M}_h), W^u_{\pm}(\mathcal{M}_h)$ .

Stable Manifold (orbits move toward the periodic orbit)



Unstable Manifold (orbits move away from the periodic orbit)

# **Transit and Nontransit Orbits**

Cylinders separate transit from nontransit orbits.
Define mappings between "bounding spheres" on either side of the potential barrier.



# **Tubes in the 3-Body Problem**

#### □ Stable and unstable manifold tubes

• Control transport through the potential barrier.



## Flux

### **Tubes of transit orbits are the relevant objects to study**

- $\Box$  Tubes determine the flux between regions  $F_{i,j}(t)$ .
- □ Note, net flux is zero for volume-preserving motion, so we consider the one-way flux.
  - Example:  $F_{J,S}(t)$  = volume of trajectories that escape from the Jupiter region into the Sun region per unit time.

#### More exotic transport between regions

Look at the intersections between the interior of stable and unstable tubes on the same energy surface.
Could be from different potential barrier saddles.



• Example: Comet transport between outside and inside of Jupiter



Look at Poincaré section intersected by both tubes.
 Choosing surface {x = constant; p<sub>x</sub> < 0}, we look at the canonical plane (y, p<sub>y</sub>).



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Relative canonical area gives relative volume of orbits.
Under certain ergodic assumptions, the relative volume can be interpreted as the probability of transition.



#### Canonical Plane $(y, p_y)$

# Mixing

By keeping track of the intersections of the tubes, one can describe the mixing of different regions  $(T_{i,j}(t))$ .

• It can get messy fast!



(from Jaffé, Farrelly and Uzer [1999])

# **Some Challenges**

□ Computationally very challenging

 $\Box$  How to handle non-transversal intersections



### **Extend to** $N \ge 3$ degrees of freedom

- □ Near equilibrium point, suppose linearized Hamiltonian vector field has eigenvalues  $\pm i\omega_j, j = 1, \ldots, N - 1$ , and  $\pm \lambda$ .
- Assume the complexification is diagonalizable.Hamiltonian normal form theory tranforms

Hamiltonian into a lowest order form:

$$H(q, p) = \sum_{i=1}^{N-1} \frac{\omega_i}{2} \left( p_i^2 + q_i^2 \right) + \lambda q_N p_N.$$

 $\Box$  Equilibrium point is of type center  $\times \cdots \times$  center  $\times$  saddle (N - 1 centers).

#### Multidimensional "saddle point"

□ For fixed energy H = h, energy surface  $\simeq S^{2N-2} \times \mathbb{R}$ . □ Constants of motion:  $I_j = q_j^2 + p_j^2, j = 1, \dots, N-1$ , and  $I_N = q_N p_N$ .



#### The N Canonical Planes

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 $\Box \text{ Normally hyperbolic invariant manifold}$ at  $q_N = p_N = 0$ ,

$$\mathcal{M}_{h} = \sum_{i=1}^{n-1} \frac{\omega_{i}}{2} \left( p_{i}^{2} + q_{i}^{2} \right) = h > 0.$$

Note that  $\mathcal{M}_h \simeq S^{2N-3}$ , not a single trajectory.

□ Four "cylinders" of asymptotic orbits: the stable and unstable manifolds  $W^s_{\pm}(\mathcal{M}_h), W^u_{\pm}(\mathcal{M}_h)$ , which have the structure  $S^{2N-3} \times \mathbb{R}$ .

- □ Transport between regions is mediated by the "higher dimensional tubes"
- □ Compute fluxes, transition probabilities, etc.



#### • Example: restricted three-body problem (3D)



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# **Future Directions**

#### **Future Directions**

• Compute fluxes, transition probabilities in 2 and 3 degree of freedom systems

#### • Determine statistical laws

- For one energy
- Over a range of energies
- Is ergodic assumption valid?
- Equilibrium distribution?
- Relaxation time to equilibrium?

#### • Apply to astronomical and chemical systems

- Astronomy: Compute asteroid collision probabilities, "equilibrium" distribution of asteroids and comets
- Chemistry: Compute reaction rates

#### • Combine with control

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